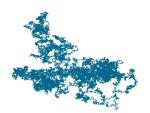
Ergodicity and convergence to equilibrium for Langevin dynamics with general potentials

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- ► *Main technical issues*: System is *degenerately damped*; randomness is also *degenerate*. Types of potential functions can make arguments harder (nonsingular vs singular).
- ► Interfaces statistical mechanics, MCMC, geometry and Boltzmann.

Markov chain $\{X_n\}$ on a finite state space $\{1,2,\ldots,d\}$ with transition P. Suppose there exists $\epsilon \in (0,1)$ and a probability η on $\{1,2,\ldots,d\}$ such that

$$P(x, A) \ge \epsilon \eta(A)$$

for all $x \in \{1, 2, \dots, d\}$, $A \subset \{1, 2, \dots, d\}$. Then for all $k \ge 1$

$$||P^k(x,\cdot) - P^k(y,\cdot)||_{TV} \le (1-\epsilon)^k ||\delta_x - \delta_y||_{TV}.$$

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Proof. Let μ_1, μ_2 be probability measures on $\{1, 2, \dots, d\}$. Define $Q(x, \cdot) = \frac{1}{1-\epsilon}P(x, \cdot) - \frac{\epsilon}{1-\epsilon}\eta(\cdot)$. Then $\mu_1P - \mu_2P = (1-\epsilon)(\mu_1Q - \mu_2Q)$

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Proof. Let μ_1, μ_2 be probability measures on $\{1, 2, \dots, d\}$. Define $Q(x, \cdot) = \frac{1}{1-\epsilon}P(x, \cdot) - \frac{\epsilon}{1-\epsilon}\eta(\cdot)$. Then $\mu_1P - \mu_2P = (1-\epsilon)(\mu_1Q - \mu_2Q)$ $\implies \|\mu_1P - \mu_2P\|_{TV} < (1-\epsilon)\|\mu_1 - \mu_2\|_{TV}.$

Let $\{X_n\}$ be a Markov chain on state space \mathcal{X} .

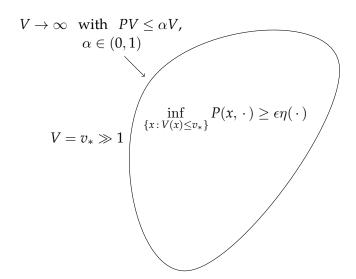
(DC) There exists $\epsilon \in (0,1)$ and a probability measure η on $\mathcal X$ such that

$$P(x, A) \ge \epsilon \eta(A)$$

for all $x \in \mathcal{X}$, $A \subset \mathcal{X}$ measurable.

- ▶ If $\mathcal{X} = \{1, 2, ..., d\}$ and $\{X_n\}$ irreducible and aperiodic, then (DC) follows for $\mathcal{P} := P^k$ (e.g. take $\nu(A) = \delta_1(A)$).
- ► If (DC) is not satisfied globally, need return times to a "small" set where (DC) is true to have exponential moments (i.e. it takes log time on average to return to small set). Use of Lyapunov structure.

CONVERGENCE PICTURE ¹



¹Harris '54; Hasminskii '80; Meyn, Tweedie '92/'93; Hairer-Mattingly '08 ∽ ∘ ∘

EXAMPLE

Stochastic gradient dynamics on \mathbf{R}^d :

$$dq_t = -\nabla U(q_t) dt + \sqrt{2} dB_t.$$

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- ▶ $U \in C^{\infty}(\mathbf{R}^d; [0, \infty))$ satisfies:
 - $U(x) \to \infty \text{ as } |x| \to \infty;$
 - $\blacktriangleright \Delta U |\nabla U|^2 \le -cU + d$ for some constants c, d > 0;

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 - $ightharpoonup \Delta U |\nabla U|^2 \le -cU + d$ for some constants c, d > 0;
 - $ightharpoonup \int e^{-U(x)} dx < \infty.$

Lyapunov structure: We have

$$LU = -|\nabla U|^2 + \Delta U \le -cU + d \implies P^t U \le e^{-ct}U + \frac{d}{c}.$$

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Doeblin Condition:

- ► Fundamental solutions of the Kolmogorov equations $(\partial_t \pm L)p = 0$, $(\partial_t \pm L^*)p = 0$ are smooth and strictly positive on $(0, \infty) \times \mathbf{R}^d$;
- ► Transition density $p_t(q, q')$ is smooth and strictly positive on $(0, \infty) \times \mathbf{R}^d \times \mathbf{R}^d$.
- ▶ (DC) follows using Lebesgue measure on a bounded set.
- ► The ϵ in (DC) is typically existential \implies quantitative minorization?²



²J. Evans '18.

POINCARÉ

Stochastic gradient dynamics on \mathbf{R}^d :

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The process q_t has a unique stationary distribution μ given by

$$\mu(dq) \propto e^{-U(q)} dq.$$

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$$\begin{split} \frac{1}{2} \frac{d}{dt} \| P^t \varphi \|_{L^2(\mu)}^2 &= \langle \frac{d}{dt} P^t \varphi, P^t \varphi \rangle = \langle L P^t \varphi, P^t \varphi \rangle \\ &= \frac{1}{2} \mu (L(P^t \varphi)^2) - \| \nabla P^t \varphi \|_{L^2(\mu)}^2 \\ &= - \| \nabla P^t \varphi \|_{L^2(\mu)}^2. \end{split}$$

Note that μ satisfies a *Poincaré inequality*. That is, there exists $\rho > 0$ such that for all $\varphi \in H^1(\mu)$ with $\mu(\varphi) = 0$ we have

$$\|\nabla\varphi\|_{L^2(\mu)}^2 \ge \rho \|\varphi\|_{L^2(\mu)}^2.$$

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Combining the above with the formal calculation gives:

$$\tfrac{1}{2} \tfrac{d}{dt} \| P^t \varphi \|_{L^2(\mu)}^2 = - \| \nabla P^t \varphi \|_{L^2(\mu)}^2 \le - \rho \| P^t \varphi \|_{L^2(\mu)}^2.$$

³Talay '00; Eckmann, Hairer '03; Hérau, Nier '04; Helffer, Nier £05 = 2000

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Hence

$$||P^t \varphi||_{L^2(\mu)} \le e^{-\rho t} ||\varphi||_{L^2(\mu)}.$$

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► Proofs of Poincaré inequality often use bounds like $\Delta U - |\nabla U|^2 < -cU + d$.

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Consider the following SDE for $x_t = (q_t, p_t)$ on $\mathcal{X} \subset \mathbf{R}^d \times \mathbf{R}^d$:

$$dq_t = p_t dt$$

$$dp_t = -\gamma p_t dt - \nabla U(q_t) dt + \sqrt{2\gamma} dB_t.$$

- ▶ B_t is a standard d-dimensional Brownian motion, $\mathcal{X} = \{U(q) < \infty\} \times \mathbf{R}^d$, $\gamma > 0$ is the friction coefficient;
- ▶ $U \in C^{\infty}(\mathcal{X}; [0, \infty))$ satisfies
 - \blacktriangleright $|\nabla U| \to \infty$ as $U \to \infty$;
 - $|\nabla^2 \dot{U}| \le \epsilon |\nabla U|^2 + C_{\epsilon}.$
- ► Hamiltonian $H(q, p) = |p|^2/2 + U(q)$ with stationary distribution ν on \mathcal{X}

$$\nu(dqdp) \propto e^{-H(q,p)} dq dp = e^{-\frac{|p|^2}{2}} e^{-U(q)} dq dp.$$

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Lyapunov?

- ► Generator: $\mathcal{L} = p \cdot \nabla_q \gamma p \cdot \nabla_p \nabla U(q) \cdot \nabla_p + \gamma \Delta_p$.
- $\blacktriangleright \mathcal{L}H(q,p) = -\gamma |p|^2 + \gamma d.$

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Conclusion: Langevin dynamics is not pointwise contractive.

AVERAGING

Example:
$$d = 1$$
, $\gamma = 1$, $U(q) = \frac{|q|^4}{4} + \frac{1}{2|q|^2}$

MORE PARTICLES

Example:
$$d = 3$$
, $\gamma = 1$, $U(q) = \sum_{i} |q_{i}|^{2} + \sum_{i \neq j} |q_{i} - q_{j}|^{-1.3}$

AVERAGING

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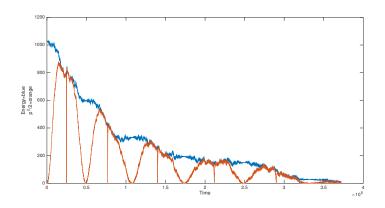


Figure: $p(t)^2/2$ and H(q(t), p(t)) plotted over $t \in [0, 4]$.

AVERAGING: HOW TO LEVERAGE?

Lyapunov: Let Av(f)(q, p) be the average value of f along Hamiltonian orbit containing (q, p), and

$$\mathcal{H}=p\cdot\nabla_q-\nabla U\cdot\nabla_p.$$

Then

$$\int_0^t |p_s|^2 ds = t \operatorname{Av}(|P|^2)(q, p) + \int_0^t |p_s|^2 - \operatorname{Av}(|P|^2)(q, p) ds.$$

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Use $V = H + \psi$ where ψ is lower-order and "satisfies"

$$\mathcal{H}\psi = |p|^2 - \operatorname{Av}(|P|^2)(q, p).$$

THE "pq TRICK"

For d = 1 and $U(q) = q^{2n}/2n$,

$$\mathcal{H}(pq) = (1+n)p^2 - 2nH(q,p) = (1+n)p^2 - \frac{1}{n+1}Av(P^2)(q,p).$$

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Polynomial-like potentials ($|\nabla^2 U| \le C |\nabla U|^1 + D$):

- ► D.Talay '00;
- ► L. Wu ′01;
- ► Mattingly/Stuart/Higham '02;
- ► Rey-Bellet '06;
- ► Zimmer '17 and Eberle, Guillin, Zimmer '19.

Different choices of ψ and general potentials (1 \mapsto 2):

- ► Cooke, H, Mattingly, McKinley, Schmidler '17;
- ► H, Mattingly '19;
- ► Lu, Mattingly '20.



POINCARÉ?

Recall for $\varphi \in L^2(\nu)$ with $\nu(\varphi) = 0$:

$$\|\mathcal{P}^t \varphi\|_{L^2(\nu)}^2 - \|\varphi\|_{L^2(\nu)}^2 = -2\gamma \int_0^t \|\nabla_p \mathcal{P}^s \varphi\|_{L^2(\nu)}^2 ds,$$

so we hope that

$$\int_0^t \|\mathcal{P}^s \varphi\|_{L^2(\nu)}^2 \, ds \lesssim \int_0^t \|\nabla_p \mathcal{P}^s \varphi\|_{L^2(\nu)}^2 \, ds.$$

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Idea 1: Follow the flow: If $\varphi_t = \mathcal{P}^t \varphi$ and $\|\cdot\| = \|\cdot\|_{L^2(\nu)}$, then

$$\frac{1}{2} \frac{d}{dt} \|\nabla_p \varphi_t\|^2 = \langle \nabla_p \mathcal{L} \varphi_t, \nabla_p \varphi_t \rangle
= \langle [\nabla_p, \mathcal{L}] \varphi_t, \nabla_p \varphi_t \rangle + \langle \mathcal{L} \nabla_p \varphi_t, \nabla_p \varphi_t \rangle
= \langle (\nabla_q - \gamma \nabla_p) \varphi_t, \nabla_p \varphi_t \rangle - \gamma \|\nabla_p^2 \varphi_t\|^2.$$

AVERAGING: HOW TO LEVERAGE THIS?

Idea 1: Follow the flow: If $\varphi_t = \mathcal{P}^t \varphi$ and $\tilde{\nabla} := \nabla_q - \gamma \nabla_p$, then

$$\begin{split} \frac{1}{2} \frac{d}{dt} \|\tilde{\nabla}\varphi_t\|_{L^2(\nu)}^2 &= \langle [\tilde{\nabla}, \mathcal{L}]\varphi_t, \tilde{\nabla}\varphi_t \rangle + \langle \mathcal{L}\tilde{\nabla}\varphi_t, \tilde{\nabla}\varphi_t \rangle \\ &= -\gamma \|\tilde{\nabla}\varphi_t\|^2 - \gamma \|\nabla_p \tilde{\nabla}\varphi_t\|^2 - \langle \nabla^2 U \nabla_p \varphi_t, \tilde{\nabla}\varphi_t \rangle \end{split}$$

AVERAGING: HOW TO LEVERAGE THIS?

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Hence use the modified $H^1(\nu)$ norm

$$|||\varphi|||^2 := c_1 ||\varphi||^2_{L^2(\nu)} + c_2 ||\nabla_p \varphi||^2 + c_3 ||\tilde{\nabla} \varphi||^2.$$

- ▶ Desvillettes, Villani '01; Hérau, Nier '04; Hellfer, Nier '05; Mouhot, Neumann '06, Hérau '07;
- ► Villani ′09;
- ▶ Conrad and Grothaus '10; Grothaus and Stilgenbauer '15;
- ▶ Baudoin '17, Monmarché '19;
- ► Cattiaux, Guillin, Monmarché, Zhang '17 and Baudoin, Gordina, H '21.



Idea 2: Construct a norm equivalent to $L^2(\nu)$ *instead:*

$$\|\varphi\|_{1+\delta A}^2 := \|\varphi\|_{L^2(\nu)}^2 + \delta \langle A\varphi, \varphi \rangle.$$

If $\varphi_t = \mathcal{P}^t \varphi$ and $\nu(\varphi) = 0$, then

$$\frac{d}{dt}\langle A\varphi_t, \varphi_t \rangle = \langle (\mathcal{L}^{\dagger}A + A\mathcal{L})\varphi_t, \varphi_t \rangle
= \langle A\mathcal{H}\Pi\varphi_t, \varphi_t \rangle + R(\varphi_t)$$

where

$$\Pi\varphi(q) = \frac{1}{\sqrt{(2\pi)^d}} \int_{\mathbf{R}^d} \varphi(q, p) e^{-\frac{|p|^2}{2}} dp.$$

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$$\frac{d}{dt}\langle A\varphi_t, \varphi_t \rangle = \langle A\mathcal{H}\Pi\varphi_t, \varphi_t \rangle + R(\varphi_t),$$

so pick $A = -(\mathcal{H}\Pi)^{\dagger}$ so that

$$\langle A\mathcal{H}\Pi\varphi_t, \varphi_t \rangle = -\|\mathcal{H}\Pi\varphi_t\|^2 = -\|p \cdot \nabla_q \Pi\varphi_t\|^2 = -c\|\nabla_q \Pi\varphi_t\|^2.$$

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Note: *A* above is not bounded on $L^2(\nu)$ so need to renormalize:

$$A\varphi = -(1 + (\mathcal{H}\Pi)^{\dagger}(\mathcal{H}\Pi))^{-1}(\mathcal{H}\Pi)^{\dagger}\varphi = \mathbf{E}_q \int_0^{\infty} e^{-s}\mathcal{H}\Pi\varphi(q_s) ds$$

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- ► Hérau '06;
- ▶ Dolbeault, Mouhot, Schmeiser '09, '15;
- ► Grothaus and F-Y Wang '19;
- ► Leimkuhler, Sachs, Stoltz '20;
- ► Camrud, Gordina, H, Stoltz '21

Idea 3: Don't change the norm!

In other words, show that for $\varphi \in L^2(\nu)$ with $\nu(\varphi) = 0$:

$$\frac{1}{\tau} \int_t^{t+\tau} \|\mathcal{P}^s \varphi\|_{L^2(\nu)}^2 ds \leq \|\varphi\|_{L^2(\nu)}^2 e^{-\lambda(\tau)t}.$$

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Time-averaged Poincaré?

$$c \int_{t}^{t+\tau} \|\mathcal{P}^{s}\varphi\|_{L^{2}(\nu)}^{2} ds \leq \int_{t}^{t+\tau} \|\nabla_{p}\mathcal{P}^{s}\varphi\|_{L^{2}(\nu)}^{2} ds$$

Hörmander's condition:

Let $U \subset \mathbf{R}^d$ be open, bounded and X_0, X_1, \dots, X_r be $C^{\infty}(U)$ vector fields. We say that X_0, X_1, \dots, X_r satisfies *Hörmander's* condition on U if for every $x \in U$, the list

$$X_{j_1}(x),$$
 $j_1 = 0, 1, ..., r$
 $[X_{j_1}, X_{j_2}](x),$ $j_1, j_2 = 0, 1, ..., r$
 $[X_{j_1}[X_{j_2}, X_{j_3}]](x),$ $j_1, j_2, j_3 = 0, 1, ..., r$
 \vdots \vdots

contains a basis of \mathbf{R}^d .

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 $[X_{j_1}[X_{j_2}, X_{j_3}]](x),$ $j_1, j_2, j_3 = 0, 1, \dots, r$
 \vdots \vdots

contains a basis of \mathbf{R}^d .

Example.
$$X_0 = p\partial_q - U'(q)\partial_p - p\partial_p$$
 and $X_1 = \partial_p$. Note that $[X_1, X_0] = \partial_q - \partial_p$.



Theorem (Hörmander 1967)

Let $K \subseteq U$ and suppose $\mathcal{M} = X_0 + \sum_{j=1}^r X_j^2$ and X_0, X_1, \dots, X_r satisfies Hörmander's condition on U. Then there exists s, C > 0 such that

$$||u||_{H^s} \le C(||\mathcal{M}u||_{L^2} + ||u||_{L^2})$$

for all $u \in C_0^{\infty}(K)$.

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Actually:

$$||u||_{H^s} \le C\bigg(||u||| + |||X_0u|||'\bigg),$$

$$|||u||| := ||u||_{L^2} + \sum_{j=1}^r ||X_ju||_{L^2}, \quad |||u|||' := \sup_{|||\varphi||| \le 1} \int u\varphi \, dx$$

Example in d=1

Question: How does this help?

Example. For Langevin in d=1 with $\gamma=1$, $\varphi_t=\mathcal{P}^t\varphi$. Then

$$\partial_t \varphi_t = \mathcal{L} \varphi_t = (\mathcal{H} + \partial_p^2) \varphi_t.$$

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$$\begin{aligned} \|\varphi_{t}\|_{H^{s}} &\leq C(\|\varphi_{t}\|_{L^{2}} + \|\partial_{p}\varphi_{t}\|_{L^{2}} + \||(\partial_{t} - \mathcal{H})\varphi_{t}|||') \\ &\leq C'(\|\varphi_{t}\|_{L^{2}} + \|\partial_{p}\varphi_{t}\|_{L^{2}}). \end{aligned}$$

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Conclusion: Try to obtain Poincaré inequality of the form

$$c\int_0^\tau \|\varphi_s\|_{L^2(\nu)} ds \leq \int_0^\tau \|\nabla_p \varphi_s\|_{L^2(\nu)} ds + |||(\partial_t - \mathcal{H})\varphi_t|||_{\nu}'$$

TIME-DEPENDENT POINCARÉ INEQUALITIES

Try to obtain time-dependent Poincaré of the form:

$$c\int_0^\tau \|\varphi_s\|_{L^2(\nu)} ds \leq \int_0^\tau \|\nabla_p \varphi_s\|_{L^2(\nu)} ds + |||(\partial_t - \mathcal{H})\varphi_t|||_{\nu}'$$

- ► Y. Guo '02;
- ► Strain and Guo '04;
- ► Albritton, Armstrong, Mourrat and Novack '21;
- ► Cao, Lu and Wang '19;
- ▶ Bedrossian and Liss '21: 2D Galerkin Navier-Stokes .
- ► Brigatti '22;
- ► Brigatti and Stoltz '23.

THANK YOU!